

Plasmons and the associated nonlinearities in phosphorene nanoribbons

Exploring extreme nonlinear optical response within a novel second principles scheme

HoW exciting! 2023, Berlin

Line Jelver¹

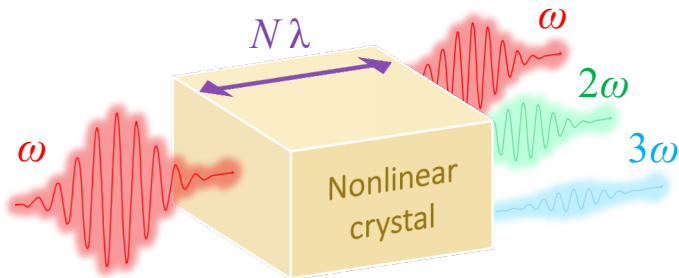
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Outline

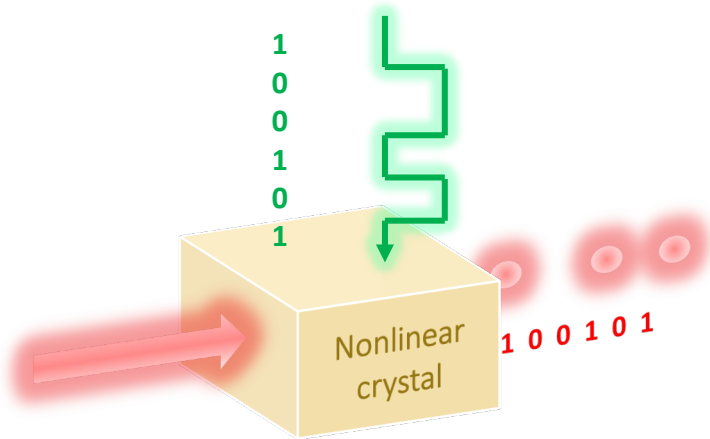
- Method
- Plasmons and the associated nonlinearities in phosphorene nanoribbons
- Electron-induced nonlinear response in atomic chains

Nonlinear optics in nanostructures

Traditional nonlinear optics

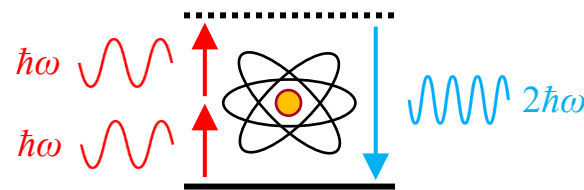
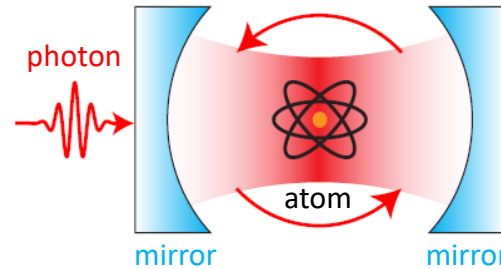


Frequency conversion

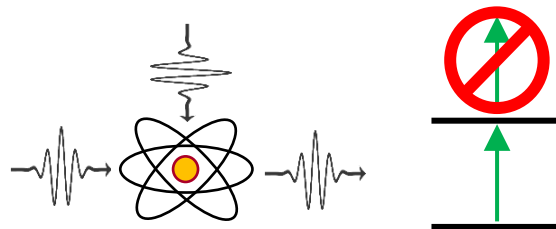


Ultrafast all-optical switching

Quantum nonlinear optics



Few-photon mixing



Single-photon blockade

Polaritons

- ✓ Hybridization of light and material excitations, such as excitons, phonons or plasmons.
- ✓ Beating the diffraction limit and enhances the field confinement which boost the light-matter interaction.

Solid-state HHG

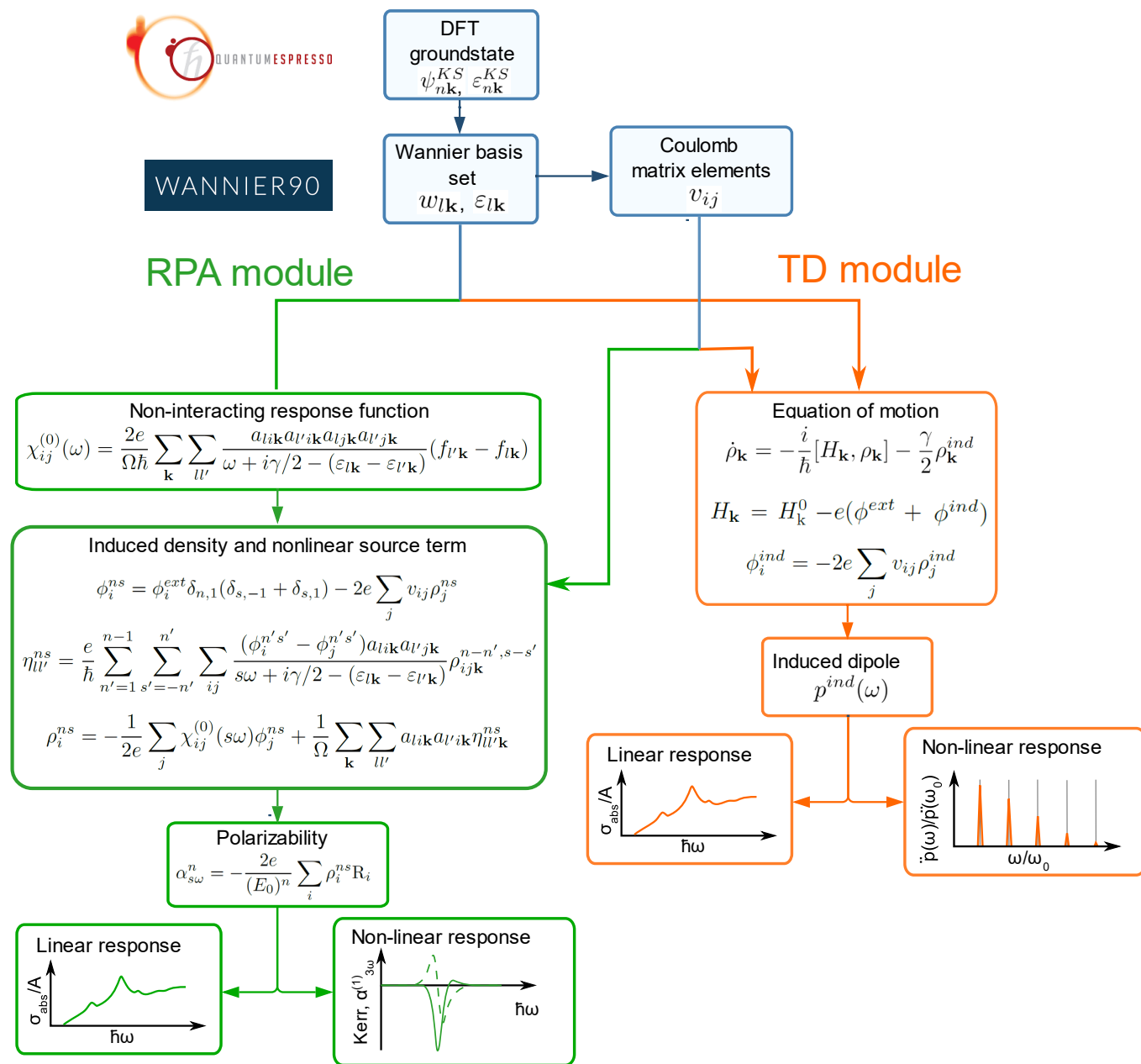
- ✓ Source of coherent high-frequency radiation which can be processed to produce **attosecond** optical pulses.

L. Jelver and J. D. Cox, arXiv:2305.09532

Method

Second principles optical response calculations

- ✓ Computationally feasible
- ✓ Offers chemical intuition
- ✓ Enables the exploration of system sizes around 10 nm's
 - ✓ The boundary between the quantum and classical regime
 - ✓ Structure sizes which enables the conversion of free light into polaritons



Perturbative expansion of the density

- Monochromatic illumination
- Expanding the density in terms of **perturbative order n** and **harmonic order s**
- The induced density can be found from the response to the induced potential

$$\rho(t) = \sum_{n,s} \rho^{ns} e^{-is\omega t}$$

Response from order (n,s)

Response from lower orders

$$\rho_l^{ns} = \sum_{l'} \chi_{ll'}^{(0)}(s\omega) \phi_{l'}^{ns} + \frac{b}{2\pi} \int_{-\pi/b}^{\pi/b} dk \sum_{jj'} a_{jl,k} a_{j'l,k}^* \eta_{jj',k}^{ns}$$



$$\alpha_{s\omega}^n = -\frac{2e}{(E_0)^n} \sum_i \rho_i^{ns} R_i$$

$$\chi_{ll'}^{(0)}(s\omega) = \frac{2e^2}{\hbar} \frac{b}{2\pi} \int_{-\pi/b}^{\pi/b} dk \sum_{jj'} (f_{j',k} - f_{j,k}) \frac{a_{jl,k} a_{j'l,k}^* a_{j'l',k} a_{j'l',k}^*}{s\omega + i/2\tau - (\epsilon_{j,k} - \epsilon_{j',k})}$$

$$\phi_l^{ns} = \mathbf{r}_l \cdot \mathbf{E}^{\text{ext}} \delta_{n,1} (\delta_{s,-1} + \delta_{s,1}) + \sum_{l'} v_{ll'} \rho_{l'}^{ns}$$

$$\eta_{jj',k}^{ns} = \frac{2e^2}{\hbar} \sum_{n'=1}^{n-1} \sum_{s'=-n'}^{n'} \sum_{ll'} \frac{(\phi_l^{n's'} - \phi_{l'}^{n's'}) a_{jl,k}^* a_{j'l',k}}{s\omega + i/2\tau - (\epsilon_{j,k} - \epsilon_{j',k})} \rho_{ll',k}^{n-n',s-s'}$$

Single-particle density matrix formulation

- Time-dependent simulations solving the equation of motion

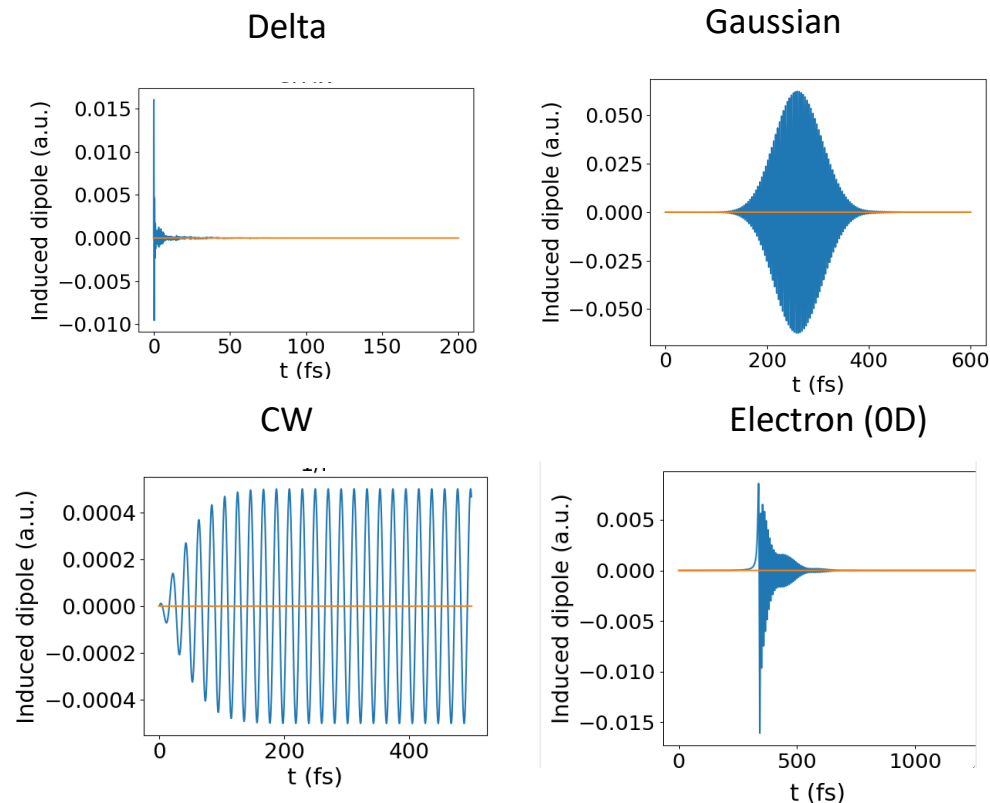
$$\frac{\partial \rho}{\partial t} = -\frac{i}{\hbar} [\mathcal{H}^{\text{TB}} - e\phi, \rho] - \frac{1}{2\tau} (\rho - \rho^{(0)}),$$

$$\phi_l = \underbrace{\mathbf{r}_l \cdot \mathbf{E}}_{\text{External potential}} + \sum_{l'} \underbrace{v_{ll'} \rho_{l'}^{\text{ind}}}_{\text{Induced potential}}$$

$$\rho_l^{\text{ind}} = -(eb/\pi) \int_{-\pi/b}^{\pi/b} dk (\rho_{ll,k} - \rho_{ll,k}^{(0)})$$

$$\mathbf{p}^{\text{ind}}(t) = \sum_l \mathbf{r}_l \rho_l^{\text{ind}}(t)$$

External potentials



L. Jelver and J. D. Cox, arXiv:2305.09532

Plasmons and the associated nonlinearities in phosphorene nanoribbons

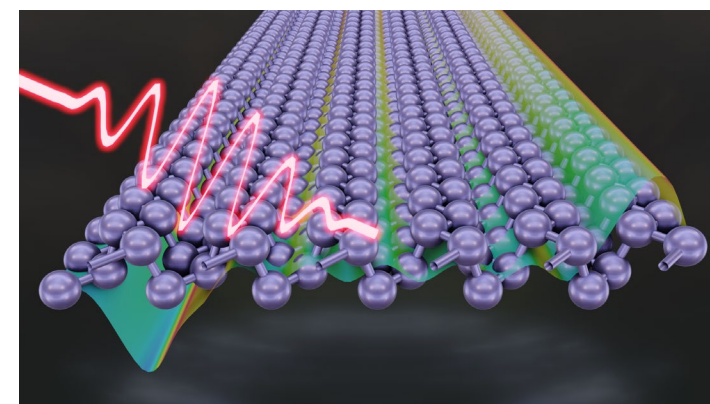
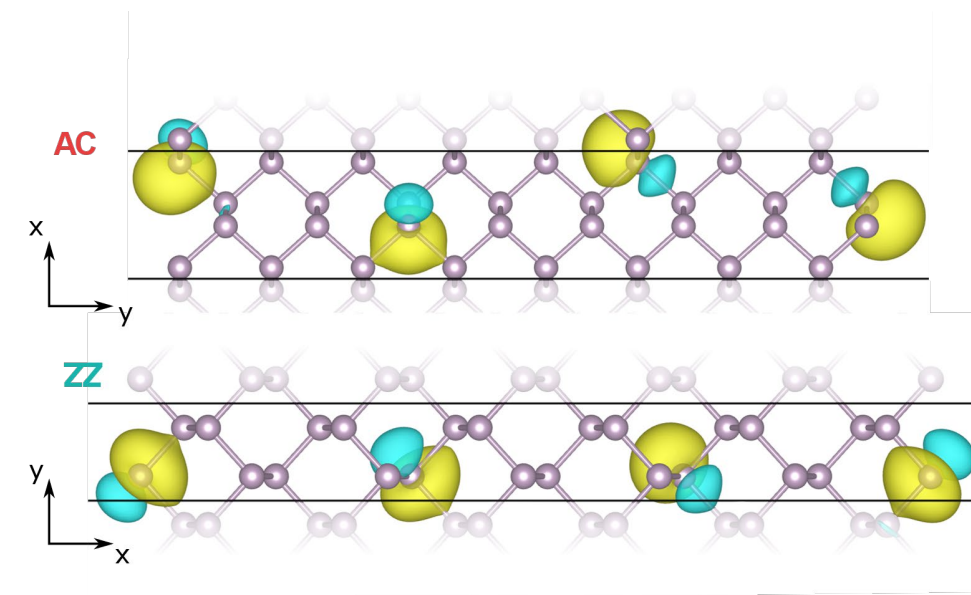
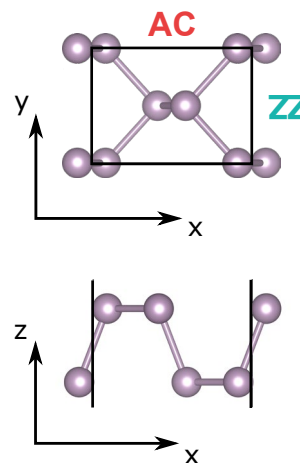
Phosphorene nanoribbons (PNRs)

2D materials

- ✓ Confinement in the out-of-plane direction
- ✓ Experience less screening which increases the probability of e-h recollisions.

Phosphorene

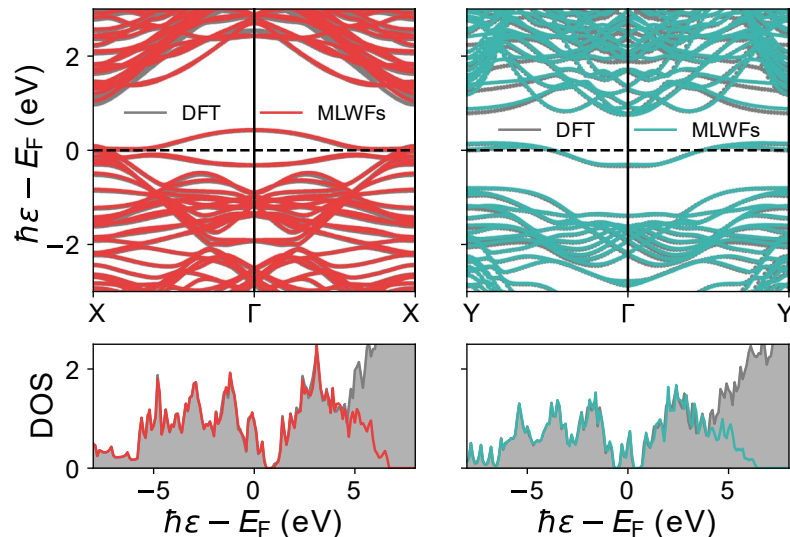
- ✓ Highly tunable anisotropic plasmonic resonances
- ✓ Strong light-matter interaction
- ✓ Extended system simulations show superior HHG compared to e.g. Graphene and MoS₂ [1]
- ✓ Can be nanostructured to break inversion symmetry allowing for even harmonics.



[1] Z.-Y. Chen and R. Qin, *Nanoscale* **11**, 16377 (2019).

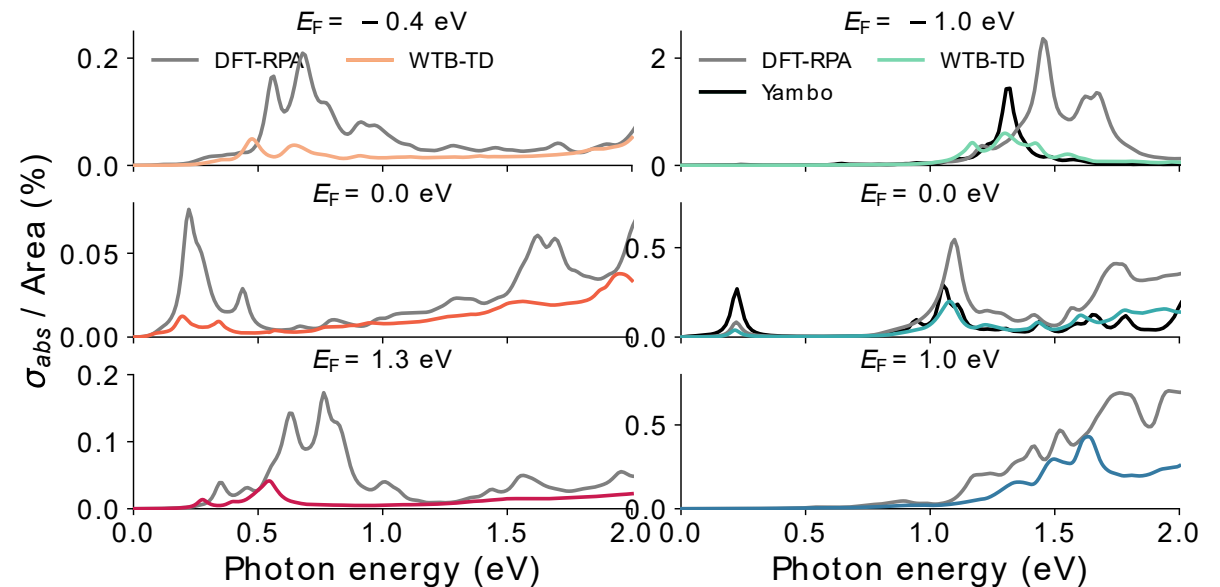
Comparison to ab-initio results

- ✓ The absorption cross section is obtained from the polarizability
- ✓ Agreement of single-particle bands 5 eV around the Fermi level.
- ✓ Qualitative features of the absorption as found from ab-initio calculations are reproduced.



GPAW!

Yambo 



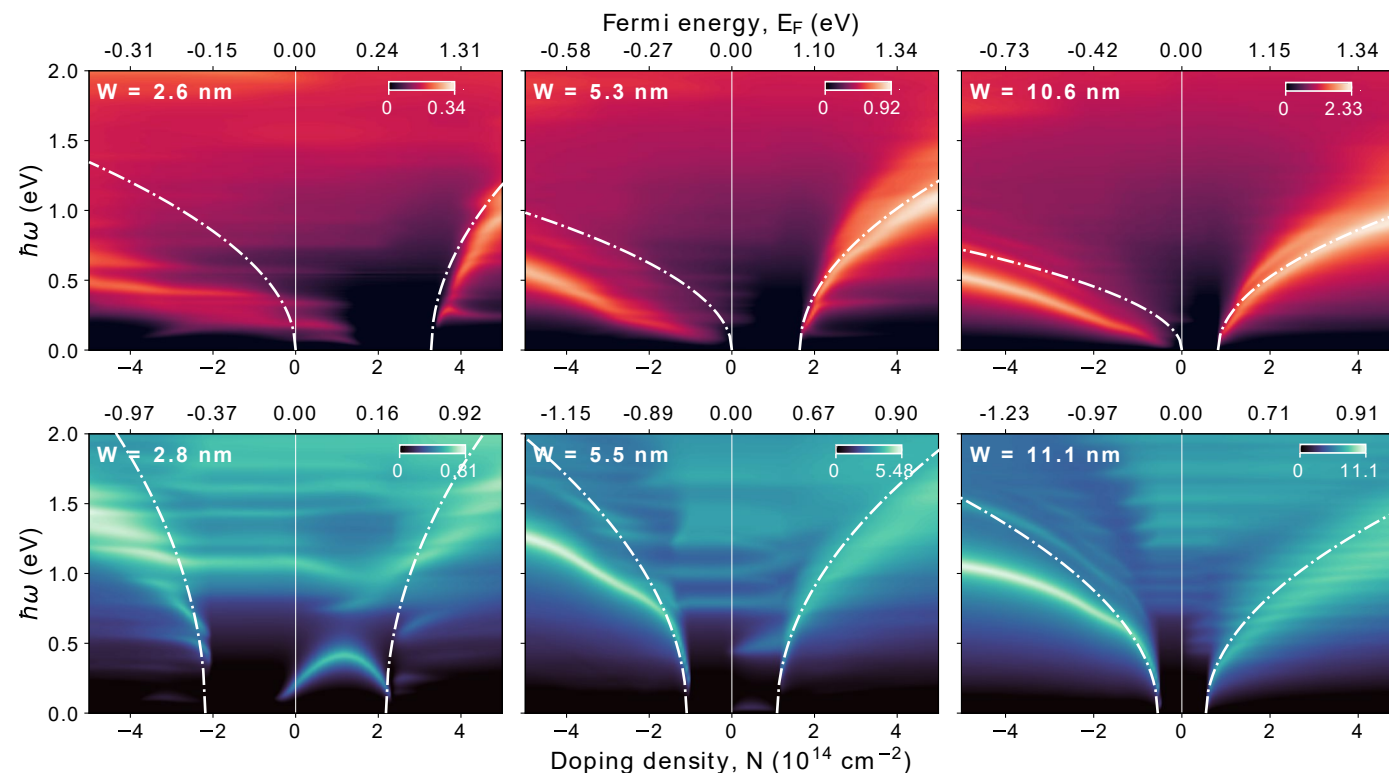
Tunability of the plasmonic modes

- ✓ Plasmonic resonances in IR and MIR range with **tunability over ≈ 1 eV**.
- ✓ Landau damping especially pronounced in hole doped ribbons.
- ✓ Agreement with analytical model of plasmonic modes in nanostructures. [1,2]

$$\omega_{\alpha}^{e/h} = \frac{1}{\sqrt{-\pi\eta_1\epsilon}} \sqrt{\frac{\mathcal{D}_{\alpha}^{e/h}}{W}},$$

$$\sigma_{\alpha}^{e/h}(\omega) = \frac{i\mathcal{D}_{\alpha}^{e/h}}{\pi(\omega + i\tau^{-1})},$$

$$\mathcal{D}_{\alpha}^{e/h} = \pi e^2 |N| / m_{\alpha,c/v}$$



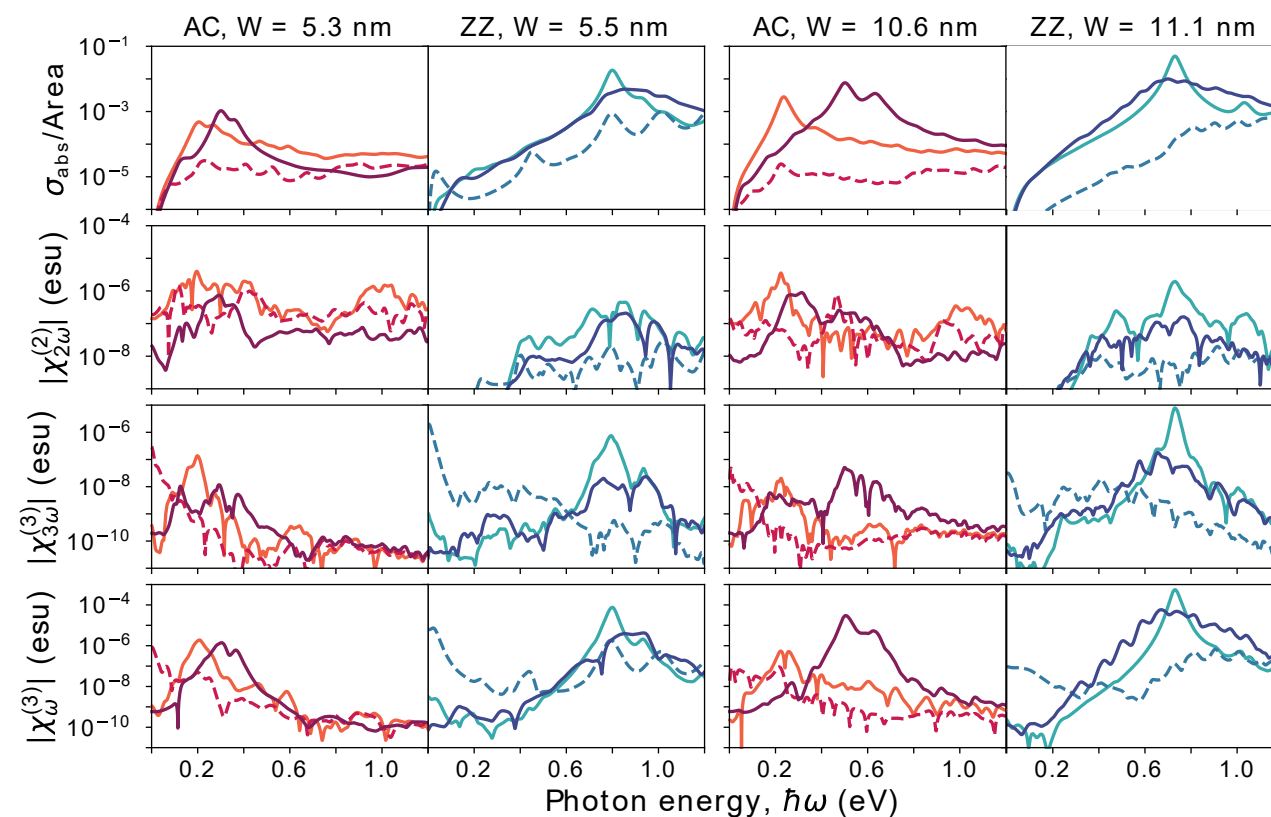
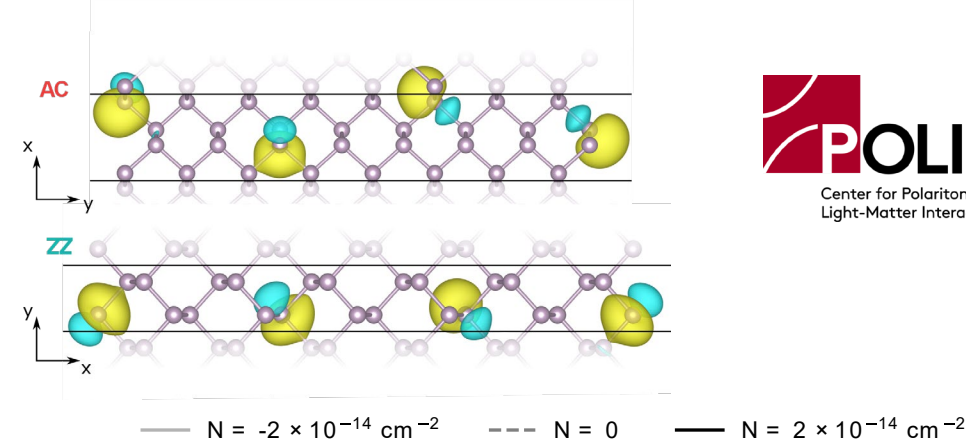
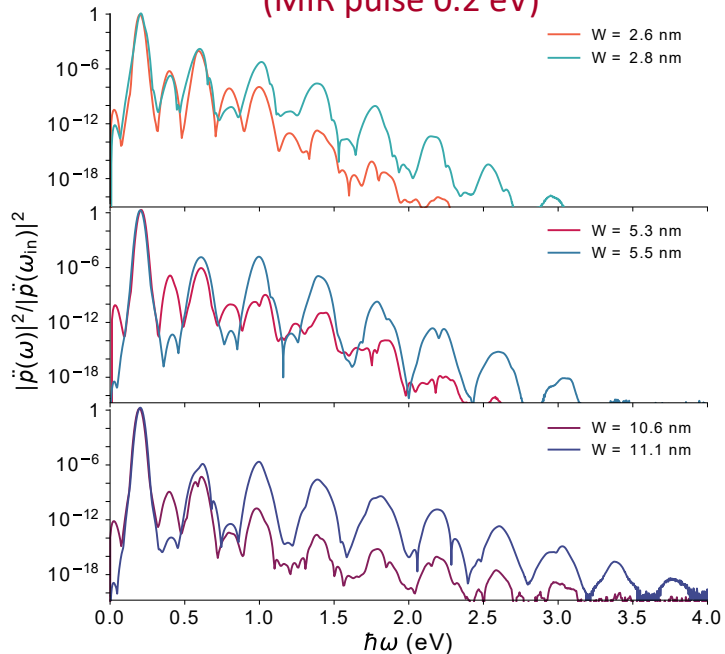
[1] R. Yu, J. D. Cox, J. R. M. Saavedra, and F. J. Garcíá De Abajo, *Analytical Modeling of Graphene Plasmons*, ACS Photonics **4**, 3106 (2017).

[2] T. Low, R. Roldán, H. Wang, F. Xia, P. Avouris, L. M. Moreno, and F. Guinea, *Plasmons and Screening in Monolayer and Multilayer Black Phosphorus*, Phys. Rev. Lett. **113**, 5 (2014).

Nonlinear response in PNRs

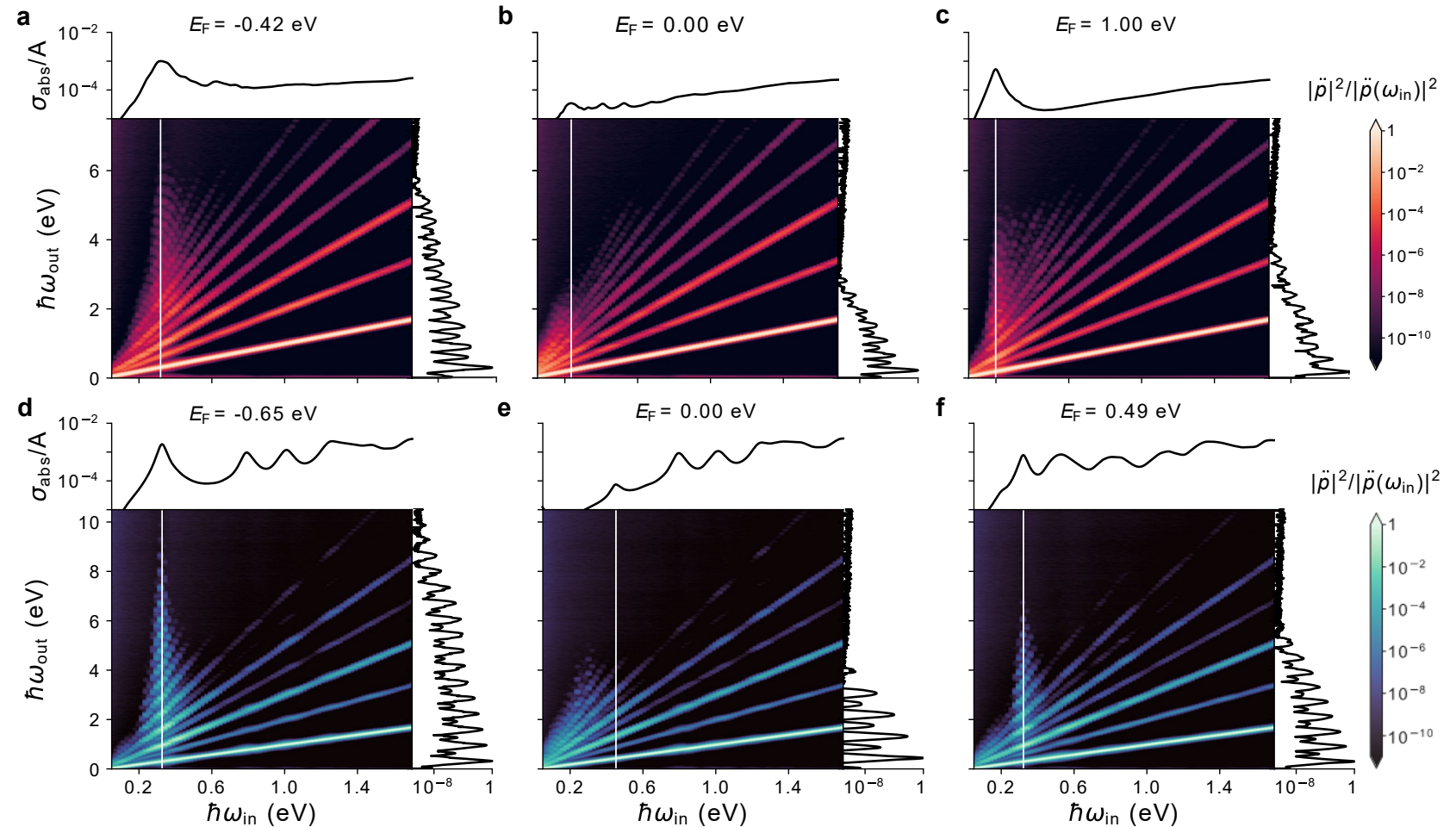
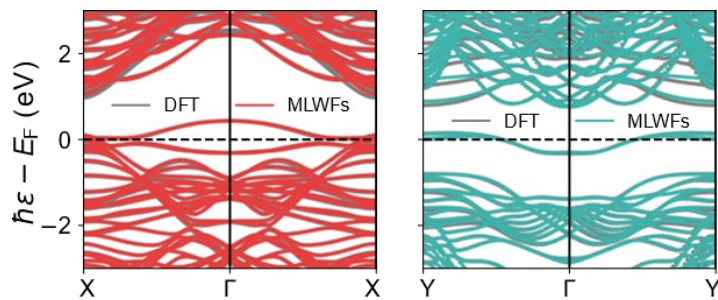
- ✓ 2nd order susceptibilities up to the order of 10^{-5} esu.
- ✓ 3rd order susceptibilities in the order of 10^{-5} esu.
- ✓ UV-range emission achieved for the two widest ribbons.

High-harmonic generation from pristine ribbons
(MIR pulse 0.2 eV)



Plasmon enhanced HHG in 5 nm ribbons

- ✓ Plasmon-assisted HHG is achieved.
- ✓ Pushes the emission towards the extreme-UV range (>10.3 eV)
- ✓ Large bandgap in hole-doped PNRs reduces quenching of the response by interband transitions.



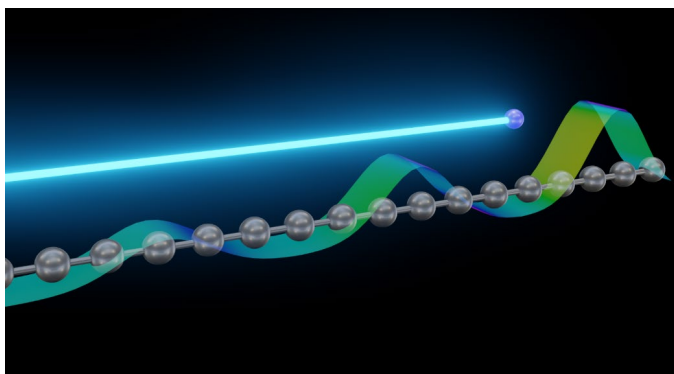
Electron-induced nonlinearity in atomic chains

L. Jelver, *et. al.*, Phys. Rev. Research 5, L022015 (2023), DOI: [10.1103/PhysRevResearch.5.L022015](https://doi.org/10.1103/PhysRevResearch.5.L022015)

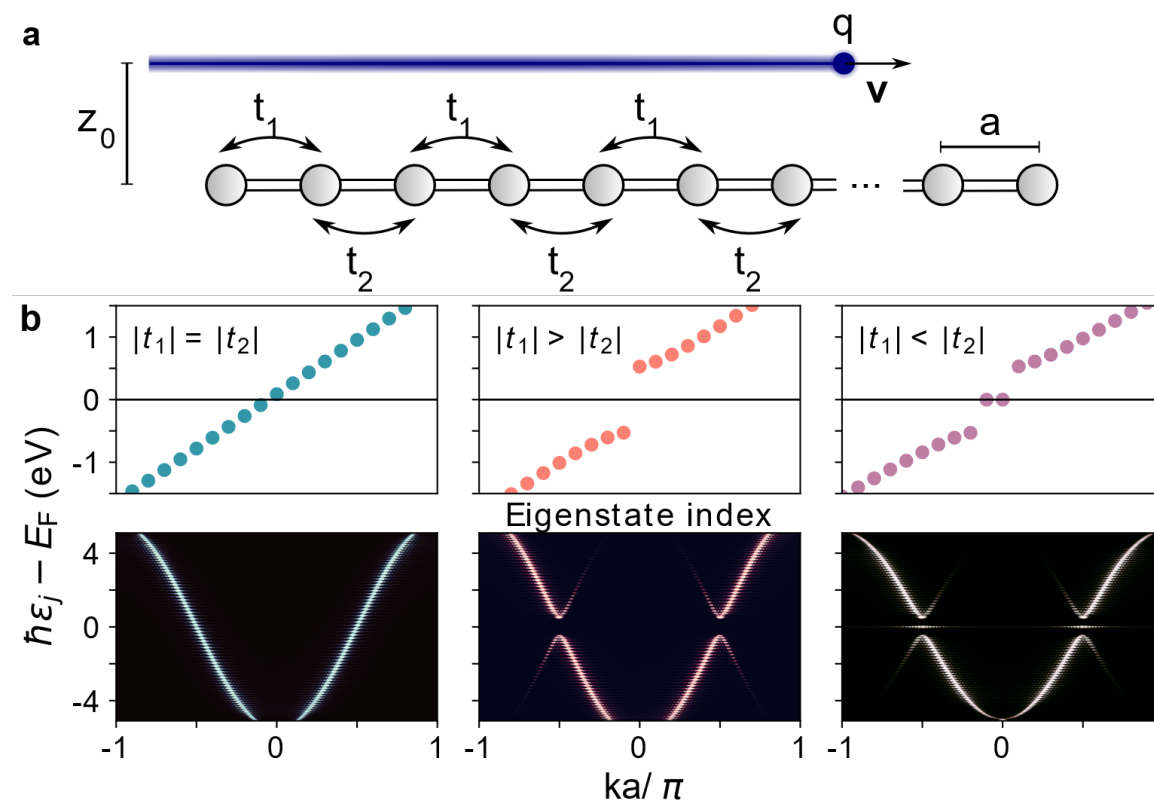
Electron-induced nonlinearity in SSH chains

- ✓ Relies on the overlap between the emanating electron field and the target.
- ✓ ~100 eV electrons can deliver optical fluences comparable to intense ultrashort optical pulses used for inducing nonlinear response.

$$\phi_l^{\text{ext}} = q |x_l \hat{\mathbf{x}} - \mathbf{r}_0 - \mathbf{v}t|^{-1}$$



$$\mathcal{H}^{\text{TB}} = t_1 \sum_{l=1}^{N/2} |2l\rangle \langle 2l-1| + t_2 \sum_{l=1}^{N/2-1} |2l+1\rangle \langle 2l| + \text{h.c.},$$



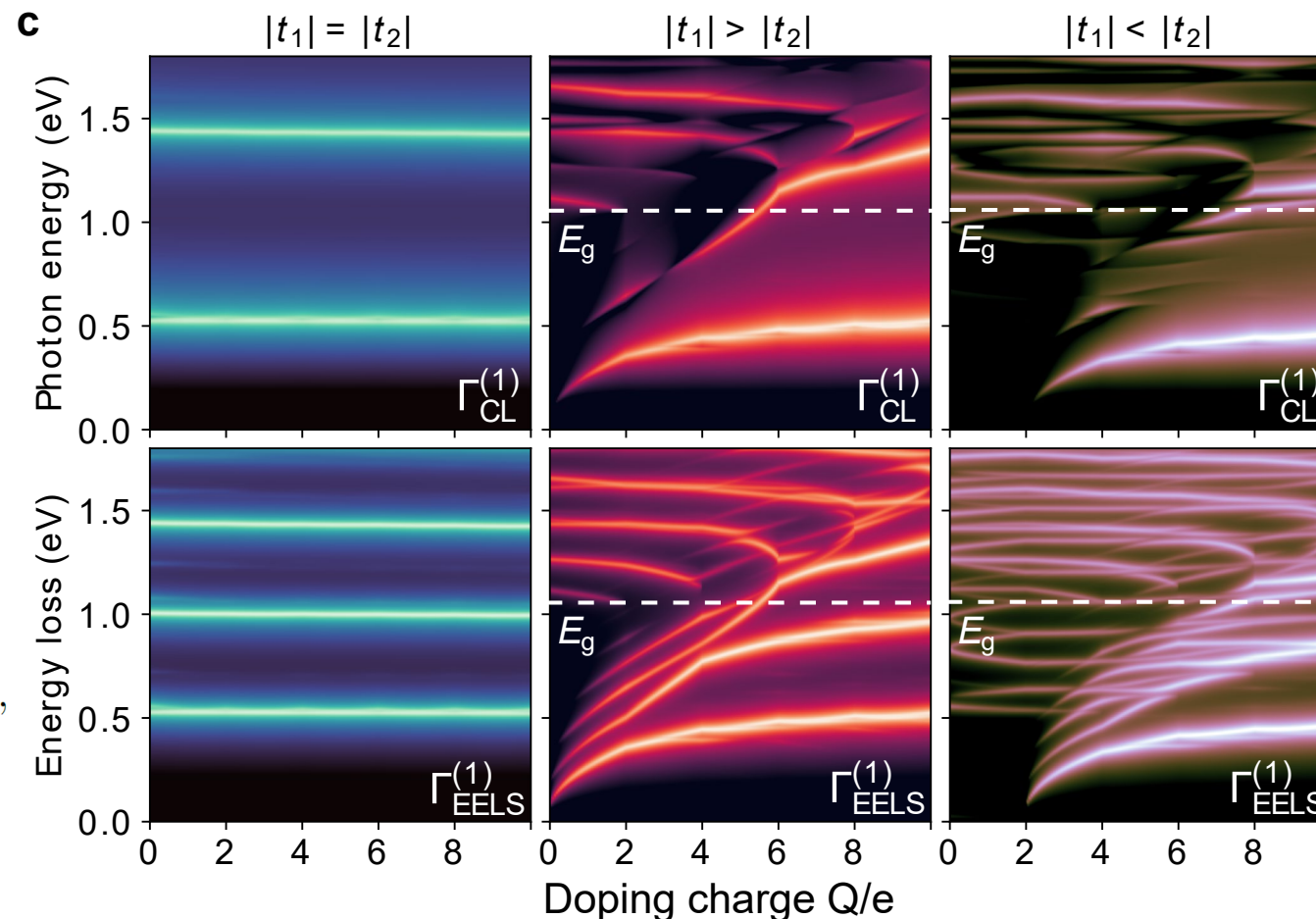
Plasmonic resonances in SSH chains

- ✓ Low-energy spectral features corresponding to plasmon resonances.
- ✓ EELS reveal the dark modes
- ✓ Tunability in the chains with electronic gaps.

$$\Gamma_{\text{CL}} = 2\omega^3 |\mathbf{p}_\omega|^2 / 3\pi\hbar c^3$$

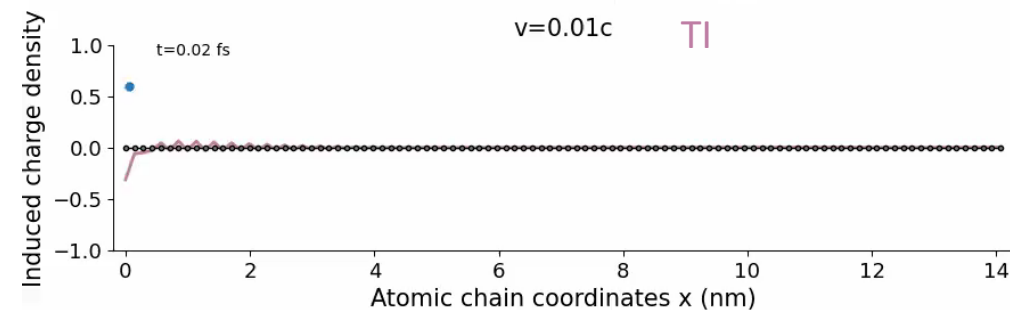
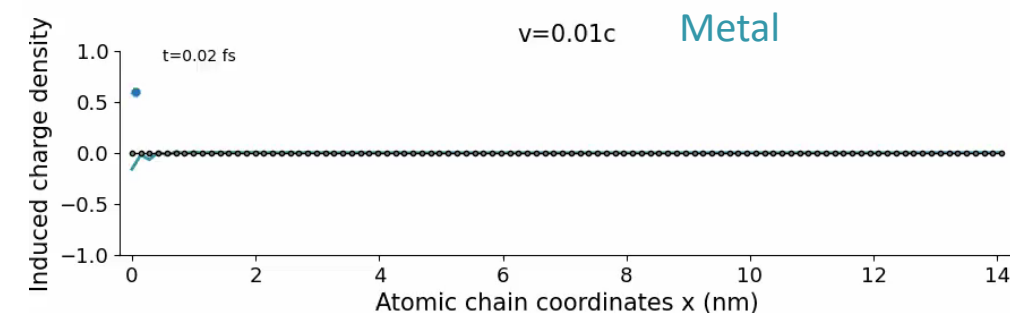
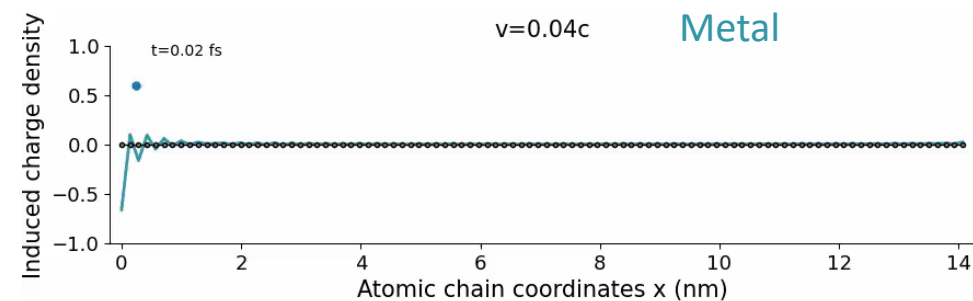
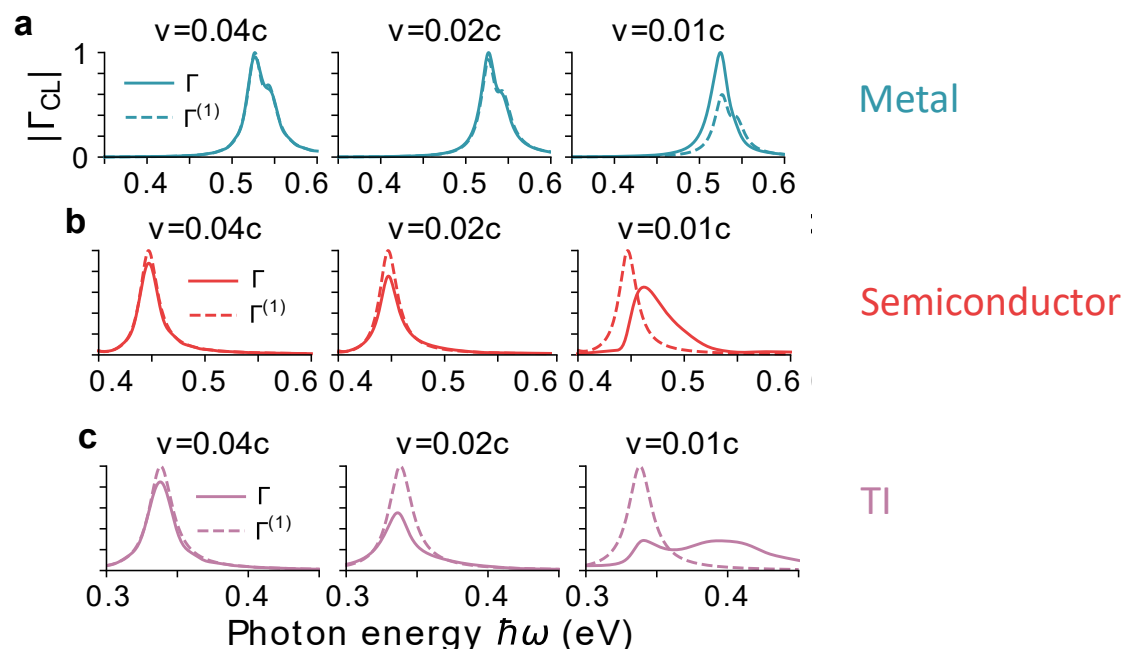
$$\Gamma_{\text{EELS}}^{\parallel} = -\frac{2q}{\pi\hbar v} \sum_l K_0 \left(\frac{\omega}{v} \sqrt{y_0^2 + z_0^2} \right) \text{Im} \left\{ \rho_{\omega,l}^{\text{ind}} e^{-i\omega x_l/v} \right\},$$

$N = 100, v = 0.01c, z_0 = 0.5 \text{ nm}$



Induced density fluctuations

- ✓ Tuning the speed to the Fermi velocity creates the most prominent **non-linear** behaviour
- ✓ The gapped chains experience the largest effect presumably due to reduced screening



Conclusions

- Novel second principles method which allows for simulating the nonlinear response of systems with up to ~ 500 atomic orbitals.

- Prospects of phosphorene nanoribbons within nano-optics
 - Excellent tunability of the plasmonic resonances
 - Tuneable and high non-linear response
 - Tuneable and large HHG yield.

- Guidelines for inducing nonlinear response by e-beams
 - Gapped systems are preferable
 - The speed of the electrons should match the Fermi velocity

Acknowledgements



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